

Blog Entry © Sunday, May 17, 2026, by James Pate Williams, Jr. Derivation of the Time Independent Free Particle Schrödinger Equation in Confocal Parabolic Coordinates

References: [Parabolic coordinates - Wikipedia](#) [finite difference method for the Laplace operator in confocal parabolic coordinates](#)

[Parabolic coordinate PDE solutions](#)

$$x = \sigma\tau \cos \varphi$$

$$y = \sigma\tau \sin \varphi$$

$$z = \frac{1}{2}(\tau^2 - \sigma^2)$$

$$dx^2 = \tau^2(\cos \varphi)^2 d\sigma^2 + \sigma^2(\cos \varphi)^2 d\tau^2 + \sigma^2\tau^2(\sin \varphi)^2 d\varphi^2$$

$$dy^2 = \tau^2(\sin \varphi)^2 d\sigma^2 + \sigma^2(\sin \varphi)^2 d\tau^2 + \sigma^2\tau^2(\cos \varphi)^2 d\varphi^2$$

$$dz^2 = \sigma^2 d\sigma^2 + \tau^2 d\tau^2$$

$$\begin{aligned} ds^2 &= \tau^2 d\sigma^2 + \sigma^2 d\tau^2 + \sigma^2\tau^2 d\varphi^2 + \sigma^2 d\sigma^2 + \tau^2 d\tau^2 \\ &= (\sigma^2 + \tau^2) d\sigma^2 + (\sigma^2 + \tau^2) d\tau^2 + \sigma^2\tau^2 d\varphi^2 \end{aligned}$$

$$g_{\sigma\sigma} = g_{\tau\tau} = \sigma^2 + \tau^2$$

$$g_{\varphi\varphi} = \sigma^2\tau^2$$

$$h_{\sigma} = h_{\tau} = \sqrt{\sigma^2 + \tau^2}$$

$$h_{\varphi} = \sigma\tau$$

$$\begin{aligned} \nabla^2 &= \frac{1}{h_{\sigma}h_{\tau}h_{\varphi}} \left[ \frac{\partial}{\partial\sigma} \left( \frac{h_{\tau}h_{\varphi}}{h_{\sigma}} \frac{\partial}{\partial\sigma} \right) + \frac{\partial}{\partial\tau} \left( \frac{h_{\sigma}h_{\varphi}}{h_{\tau}} \frac{\partial}{\partial\tau} \right) + \frac{\partial}{\partial\varphi} \left( \frac{h_{\sigma}h_{\tau}}{h_{\varphi}} \frac{\partial}{\partial\varphi} \right) \right] \\ &= \frac{1}{(\sigma^2 + \tau^2)\sigma\tau} \left[ \frac{\partial}{\partial\sigma} \left( \sigma\tau \frac{\partial}{\partial\sigma} \right) + \frac{\partial}{\partial\tau} \left( \sigma\tau \frac{\partial}{\partial\tau} \right) + \frac{\partial}{\partial\varphi} \left( \frac{\sigma^2 + \tau^2}{\sigma\tau} \frac{\partial}{\partial\varphi} \right) \right] \\ &= \frac{1}{(\sigma^2 + \tau^2)\sigma} \frac{\partial}{\partial\sigma} \left( \sigma \frac{\partial}{\partial\sigma} \right) + \frac{1}{(\sigma^2 + \tau^2)\tau} \frac{\partial}{\partial\tau} \left( \tau \frac{\partial}{\partial\tau} \right) + \frac{1}{\sigma^2\tau^2} \frac{\partial^2}{\partial\varphi^2} \end{aligned}$$

The free particle Schrödinger equation in confocal parabolic coordinates is as follows:

$$-\frac{\hbar^2}{8\pi^2m} \left\{ \frac{1}{(\sigma^2 + \tau^2)\sigma} \frac{\partial}{\partial\sigma} \left( \sigma \frac{\partial}{\partial\sigma} \right) + \frac{1}{(\sigma^2 + \tau^2)\tau} \frac{\partial}{\partial\tau} \left( \tau \frac{\partial}{\partial\tau} \right) + \frac{1}{\sigma^2\tau^2} \frac{\partial^2}{\partial\varphi^2} \right\} \psi(\sigma, \tau, \varphi) = 0$$

We now use separation of variables:

$$\psi(\sigma, \tau, \varphi) = \Sigma(\sigma, \tau)\Phi(\varphi)$$

$$\begin{aligned}
& -\frac{\hbar^2}{8\pi^2 m} \left\{ \frac{1}{(\sigma^2 + \tau^2)\sigma} \frac{\partial}{\partial \sigma} \left( \sigma \frac{\partial}{\partial \sigma} \right) + \frac{1}{(\sigma^2 + \tau^2)\tau} \frac{\partial}{\partial \tau} \left( \tau \frac{\partial}{\partial \tau} \right) + \frac{1}{\sigma^2 \tau^2} \frac{\partial^2}{\partial \varphi^2} \right\} \Sigma(\sigma, \tau) \Phi(\varphi) \\
& = -\frac{\hbar^2}{8\pi^2 m} \left\{ \frac{1}{(\sigma^2 + \tau^2)\sigma} \frac{\partial}{\partial \sigma} \left( \sigma \frac{\partial \Sigma(\sigma, \tau)}{\partial \sigma} \right) + \frac{1}{(\sigma^2 + \tau^2)\tau} \frac{\partial}{\partial \tau} \left( \tau \frac{\partial \Sigma(\sigma, \tau)}{\partial \tau} \right) \right\} \Phi(\varphi) \\
& \quad - \frac{\hbar^2}{8\pi^2 m} \frac{1}{\sigma^2 \tau^2} \frac{\partial^2 \Phi(\varphi)}{\partial \varphi^2} \Sigma(\sigma, \tau)
\end{aligned}$$

Divide by the separate wavefunctions:

$$\begin{aligned}
& -\frac{\hbar^2}{8\pi^2 \mu} \frac{\sigma^2 \tau^2}{\Sigma(\sigma, \tau)} \left\{ \frac{1}{(\sigma^2 + \tau^2)\sigma} \frac{\partial}{\partial \sigma} \left( \sigma \frac{\partial \Sigma(\sigma, \tau)}{\partial \sigma} \right) + \frac{1}{(\sigma^2 + \tau^2)\tau} \frac{\partial}{\partial \tau} \left( \tau \frac{\partial \Sigma(\sigma, \tau)}{\partial \tau} \right) \right\} \\
& = \frac{\hbar^2}{8\pi^2 \mu} \frac{1}{\Phi(\varphi)} \frac{\partial^2 \Phi(\varphi)}{\partial \varphi^2} = -\lambda^2
\end{aligned}$$

We solve the easy Phi equation:

$$\frac{d^2 \Phi(\varphi)}{d\varphi^2} + \frac{8\pi^2 \mu \lambda^2}{\hbar^2} \Phi(\varphi) = 0$$

Set the Phi multiplication constant is as illustrated:

$$\alpha^2 = \frac{8\pi^2 \mu \lambda^2}{\hbar^2}$$

$$\Phi(\varphi) = A \sin(\alpha\varphi) + B \cos(\alpha\varphi)$$

$$\frac{d\Phi}{d\varphi} = \alpha A \cos(\alpha\varphi) - \alpha B \sin(\alpha\varphi)$$

$$\frac{d^2 \Phi}{d\varphi^2} = -\alpha^2 A \sin(\alpha\varphi) - \alpha^2 B \cos(\alpha\varphi)$$

The coupled equation is:

$$\begin{aligned}
& -\frac{\hbar^2 \sigma^2 \tau^2}{8\pi^2 m} \left\{ \frac{1}{(\sigma^2 + \tau^2)\sigma} \frac{\partial}{\partial \sigma} \left( \sigma \frac{\partial \Sigma(\sigma, \tau)}{\partial \sigma} \right) + \frac{1}{(\sigma^2 + \tau^2)\tau} \frac{\partial}{\partial \tau} \left( \tau \frac{\partial \Sigma(\sigma, \tau)}{\partial \tau} \right) \right\} + \lambda^2 \Sigma(\sigma, \tau) = 0 \\
& -\frac{\hbar^2}{8\pi^2 m} \left\{ \frac{\sigma \tau^2}{(\sigma^2 + \tau^2)} \left[ \sigma \frac{\partial^2 \Sigma(\sigma, \tau)}{\partial \sigma^2} + \frac{\partial \Sigma(\sigma, \tau)}{\partial \sigma} \right] + \frac{\sigma^2 \tau}{(\sigma^2 + \tau^2)} \left[ \tau \frac{\partial^2 \Sigma(\sigma, \tau)}{\partial \tau^2} + \frac{\partial \Sigma(\sigma, \tau)}{\partial \tau} \right] \right\} \\
& \quad + \lambda^2 \Sigma(\sigma, \tau) = 0 \\
& \frac{\sigma \tau^2}{(\sigma^2 + \tau^2)} \left[ \sigma \frac{\partial^2 \Sigma(\sigma, \tau)}{\partial \sigma^2} + \frac{\partial \Sigma(\sigma, \tau)}{\partial \sigma} \right] + \frac{\sigma^2 \tau}{(\sigma^2 + \tau^2)} \left[ \tau \frac{\partial^2 \Sigma(\sigma, \tau)}{\partial \tau^2} + \frac{\partial \Sigma(\sigma, \tau)}{\partial \tau} \right] - \frac{8\pi^2 \mu \lambda^2}{\hbar^2} \Sigma(\sigma, \tau) = 0 \\
& \left[ \frac{\partial^2 \Sigma(\sigma, \tau)}{\partial \sigma^2} + \frac{1}{\sigma} \frac{\partial \Sigma(\sigma, \tau)}{\partial \sigma} \right] + \left[ \frac{\partial^2 \Sigma(\sigma, \tau)}{\partial \tau^2} + \frac{1}{\tau} \frac{\partial \Sigma(\sigma, \tau)}{\partial \tau} \right] - \frac{8\pi^2 \mu \lambda^2 (\sigma^2 + \tau^2)}{\hbar^2 \sigma^2 \tau^2} \Sigma(\sigma, \tau) = 0
\end{aligned}$$

$$\frac{\partial^2 \Sigma(\sigma, \tau)}{\partial \sigma^2} + \frac{1}{\sigma} \frac{\partial \Sigma(\sigma, \tau)}{\partial \sigma} + \frac{\partial^2 \Sigma(\sigma, \tau)}{\partial \tau^2} + \frac{1}{\tau} \frac{\partial \Sigma(\sigma, \tau)}{\partial \tau} - \frac{8\pi^2 \mu \lambda^2 (\sigma^2 + \tau^2)}{h^2 \sigma^2 \tau^2} \Sigma(\sigma, \tau) = 0$$

We now compute the corresponding finite differences:

$$\frac{\partial^2 \Sigma}{\partial \sigma^2} = \frac{\Sigma_{i+1,j} - 2\Sigma_{i,j} + \Sigma_{i-1,j}}{(\Delta\sigma)^2}$$

$$\frac{\partial \Sigma}{\partial \sigma} = \frac{\Sigma_{i+1,j} - \Sigma_{i-1,j}}{2\Delta\sigma}$$

$$\frac{\partial^2 \Sigma}{\partial \tau^2} = \frac{\Sigma_{i+1,j} - 2\Sigma_{i,j} + \Sigma_{i-1,j}}{(\Delta\tau)^2}$$

$$\frac{\partial \Sigma}{\partial \tau} = \frac{\Sigma_{i+1,j} - \Sigma_{i-1,j}}{2\Delta\tau}$$

$$\begin{aligned} & \frac{\Sigma_{i+1,j} - 2\Sigma_{i,j} + \Sigma_{i-1,j}}{(\Delta\sigma)^2} + \frac{1}{\sigma} \frac{\Sigma_{i+1,j} - \Sigma_{i-1,j}}{2\Delta\sigma} + \frac{\Sigma_{i+1,j} - 2\Sigma_{i,j} + \Sigma_{i-1,j}}{(\Delta\tau)^2} + \frac{1}{\tau} \frac{\Sigma_{i+1,j} - \Sigma_{i-1,j}}{2\Delta\tau} \\ & - \frac{8\pi^2 \mu \lambda^2 (\sigma^2 + \tau^2)}{h^2 \sigma^2 \tau^2} \Sigma_{ij} = 0 \end{aligned}$$

The previous equations have singular coordinates at the origin. Now suppose asymptotically near the origin we use the following coordinate transformation:

$$\sigma = e^u, u = \log \sigma, \tau = e^v, v = \log \tau$$

Using the chain rule differentiation formula, we obtain:

$$\frac{\partial \Sigma}{\partial \sigma} = \frac{\partial \Sigma}{\partial u} \frac{\partial u}{\partial \sigma} = \frac{1}{\sigma} \frac{\partial \Sigma}{\partial u} = e^{-u} \frac{\partial \Sigma}{\partial u}$$

$$\frac{\partial^2 \Sigma}{\partial \sigma^2} = \frac{\partial}{\partial \sigma} \left( \frac{1}{\sigma} \frac{\partial \Sigma}{\partial u} \right) = -\frac{1}{\sigma^2} \frac{\partial \Sigma}{\partial u} + \frac{1}{\sigma} \frac{\partial^2 \Sigma}{\partial \sigma \partial u} = -e^{-2u} \frac{\partial \Sigma}{\partial u} + e^{-u} \frac{\partial^2 \Sigma}{\partial u^2}$$

$$e^{-u} \frac{\partial^2 \Sigma}{\partial u^2} - e^{-2u} \frac{\partial \Sigma}{\partial u} + e^{-v} \frac{\partial^2 \Sigma}{\partial v^2} - e^{-2v} \frac{\partial \Sigma}{\partial v} - \frac{8\pi^2 \mu \lambda^2 (e^{2u} + e^{2v})}{h^2 e^{2u} e^{2v}} \Sigma = 0$$

We seem to have good behavior as the coordinates  $u$  and  $v$  jointly approach infinity.

$$\lim_{\sigma, \tau \rightarrow \infty} \frac{\sigma^2 + \tau^2}{\sigma^2 \tau^2} = \lim_{\sigma, \tau \rightarrow \infty} \frac{2\sigma + 2\tau}{2\sigma\tau^2 + 2\sigma^2\tau} = \lim_{\sigma, \tau \rightarrow \infty} \frac{4}{2\tau^2 + 4\sigma\tau + 2\sigma^2 + 4\sigma\tau} = 0$$

Now assume the eigenvalue lambda is zero:

$$\frac{\partial^2 \Sigma(\sigma, \tau)}{\partial \sigma^2} + \frac{1}{\sigma} \frac{\partial \Sigma(\sigma, \tau)}{\partial \sigma} + \frac{\partial^2 \Sigma(\sigma, \tau)}{\partial \tau^2} + \frac{1}{\tau} \frac{\partial \Sigma(\sigma, \tau)}{\partial \tau} = 0$$

$$\frac{\partial^2 \Sigma(\sigma, \tau)}{\partial \sigma^2} + \frac{1}{\sigma} \frac{\partial \Sigma(\sigma, \tau)}{\partial \sigma} = -\frac{\partial^2 \Sigma(\sigma, \tau)}{\partial \tau^2} - \frac{1}{\tau} \frac{\partial \Sigma(\sigma, \tau)}{\partial \tau} = -n^2$$

$$\Sigma(\sigma, \tau) = S(\sigma)T(\tau)$$

$$\frac{1}{S} \left( \frac{d^2 S}{d\sigma^2} + \frac{1}{\sigma} \frac{dS}{d\sigma} \right) + n^2 = 0$$

$$\frac{1}{T} \left( \frac{d^2 T}{d\tau^2} + \frac{1}{\tau} \frac{dT}{d\tau} \right) - n^2 = 0$$

$$\frac{d^2 S}{d\sigma^2} + \frac{1}{\sigma} \frac{dS}{d\sigma} + n^2 S = 0$$

$$\frac{d^2 T}{d\tau^2} + \frac{1}{\tau} \frac{dT}{d\tau} - n^2 T = 0$$

$$\sigma^2 \frac{d^2 S}{d\sigma^2} + \sigma \frac{dS}{d\sigma} + n^2 \sigma^2 S = 0$$

$$\tau^2 \frac{d^2 T}{d\tau^2} + \tau \frac{dT}{d\tau} - n^2 \tau^2 T = 0$$

$$S(\sigma) = \sum_{k=0}^{\infty} a_k \sigma^k$$

$$T(\tau) = \sum_{k=0}^{\infty} b_k \tau^k$$

$$\sum_{k=0}^{\infty} [k(k-1)a_k \sigma^{k-2} + k a_k \sigma^{k-1} + n^2 \sigma^{k+2} a_k] = 0$$

$$2 \cdot 1 a_2 + 2 a_1 + n^2 \sigma^2 a_0 = 0$$

$$3 \cdot 2 a_3 + 3 a_2 + n^2 \sigma^3 a_1 = 0$$

$$4 \cdot 3 a_4 + 4 a_3 + n^2 \sigma^4 a_2 = 0$$

$$k \cdot (k-1) a_k + k a_{k-1} + n^2 \sigma^k a_{k-2} = 0$$

$$a_k = -\frac{1}{k \cdot (k-1)} (k a_{k-1} + n^2 \sigma^k a_{k-2})$$

Suppose the eigenvalue  $n$  is zero:

$$a_k = -\frac{1}{k-1} a_{k-1}$$

Let

$$a_2 = -a_1 = -1$$

$$a_3 = \frac{1}{2}$$

$$a_4 = -\frac{1}{3 \cdot 2}$$

$$a_5 = \frac{1}{4 \cdot 3 \cdot 2}$$

$$a_k = \frac{(-1)^{k+1}}{(k-1)!} \forall n \geq 2, a_1 = 1$$

$$b_k = \frac{(-1)^{k+1}}{(k-1)!} \forall n \geq 2, b_1 = 1$$